

MIT OpenCourseWare
<http://ocw.mit.edu>

2.094 Finite Element Analysis of Solids and Fluids
Spring 2008

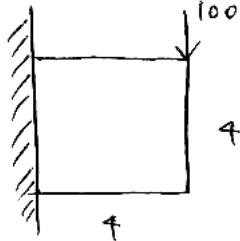
For information about citing these materials or our Terms of Use, visit: <http://ocw.mit.edu/terms>.

Lecture 6 - Finite element formulation, example, convergence

Prof. K.J. Bathe

MIT OpenCourseWare

6.1 Example

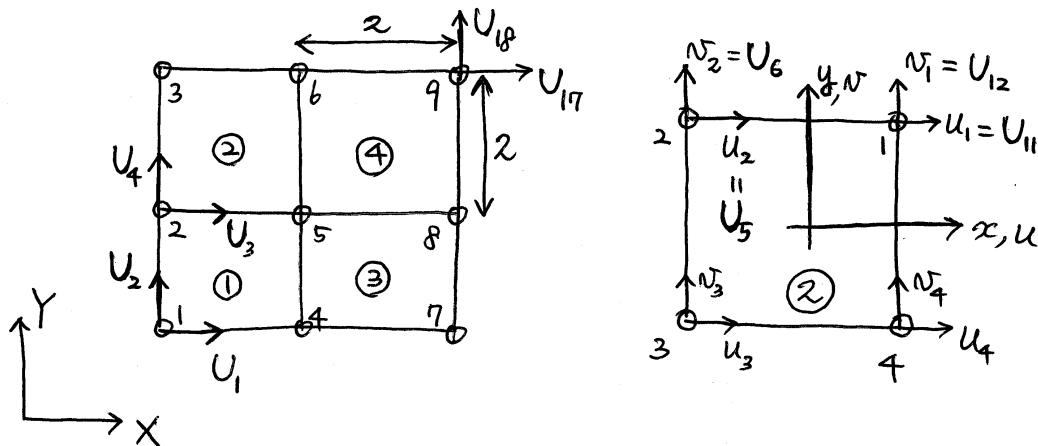
 $t = 0.1, E, \nu$ plane stressReading:
Ex. 4.6 in
the text

$$\mathbf{KU} = \mathbf{R}; \quad \mathbf{R} = \mathbf{R}_B + \mathbf{R}_s + \mathbf{R}_c + \mathbf{R}_r \quad (6.1)$$

$$\mathbf{K} = \sum_m \mathbf{K}^{(m)}; \quad \mathbf{K}^{(m)} = \int_{V^{(m)}} \mathbf{B}^{(m)T} \mathbf{C}^{(m)} \mathbf{B}^{(m)} dV^{(m)} \quad (6.2)$$

$$\mathbf{R}_B = \sum_m \mathbf{R}_B^{(m)}; \quad \mathbf{R}_B^{(m)} = \int_{V^{(m)}} \mathbf{H}^{(m)T} \mathbf{f}^{B(m)} dV^{(m)} \quad (6.3)$$

6.1.1 F.E. model



$$\mathbf{K}_{\text{el. (2)}} = \begin{bmatrix} u_1 & u_2 & u_3 & u_4 & v_1 & v_2 & v_3 & v_4 \\ \downarrow & \downarrow \\ \circ & \square & \triangle & \times & \times & \times & \times & \times \\ \vdots & & & & & & & \end{bmatrix} \quad \leftarrow u_1 \quad \vdots \quad (6.4)$$

In practice,

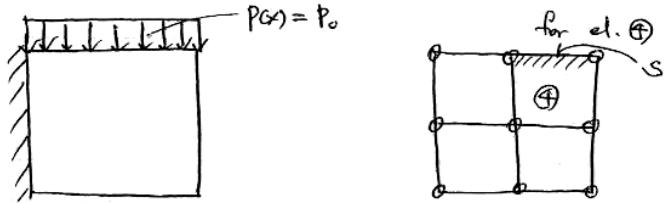
$$\mathbf{K}_{\text{el}} = \int_V \mathbf{B}^T \mathbf{C} \mathbf{B} dV; \quad \boldsymbol{\epsilon} = \mathbf{B} \begin{pmatrix} u_1 \\ \vdots \\ u_4 \\ v_1 \\ \vdots \\ v_4 \end{pmatrix} \quad (6.5)$$

where \mathbf{K} is 8x8 and \mathbf{B} is 3x8.

Assume we have \mathbf{K} (8x8) for el. (2)

$$\underbrace{\mathbf{K}_{\text{assemblage}}}_{=} = \begin{bmatrix} U_1 & U_2 & U_3 & U_4 & U_5 & \cdots & U_{11} & \cdots & U_{18} \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & & \downarrow & & \downarrow \\ \times & \times \\ \vdots & & \vdots & & \vdots & & \vdots & & \\ \vdots & & \vdots & & \vdots & & \vdots & & \\ \Delta & & \square & & \bigcirc & & & & \\ \vdots & & \vdots & & \vdots & & & & \\ \vdots & & \vdots & & \vdots & & & & \\ \times & \times \end{bmatrix} \begin{array}{l} \leftarrow U_1 \\ \vdots \\ \leftarrow U_{11} \\ \vdots \\ \leftarrow U_{18} \end{array} \quad (6.6)$$

Consider,



$$\mathbf{R}_S = \int_S \mathbf{H}^{S^T} \mathbf{f}^S dS; \quad \mathbf{H}^S = \mathbf{H} \Big|_{\text{on surface}} \quad (6.7)$$

$$\mathbf{H} = \begin{bmatrix} h_1 & h_2 & h_3 & h_4 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & h_1 & h_2 & h_3 & h_4 \end{bmatrix} \begin{array}{l} \leftarrow u(x, y) \\ \leftarrow v(x, y) \end{array} \quad (6.8)$$

$$\mathbf{u} = \begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ u_4 \\ v_1 \\ \vdots \\ v_4 \end{pmatrix} \quad (6.9)$$

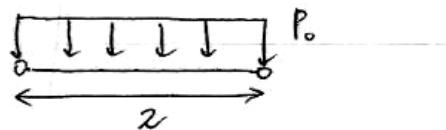
$$\mathbf{H}^S = \mathbf{H} \Big|_{y=+1} \quad (6.10)$$

$$= \begin{bmatrix} \frac{1}{2}(1+x) & \frac{1}{2}(1-x) & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{2}(1+x) & \frac{1}{2}(1-x) & 0 & 0 \end{bmatrix} \quad (6.11)$$

From (6.7);

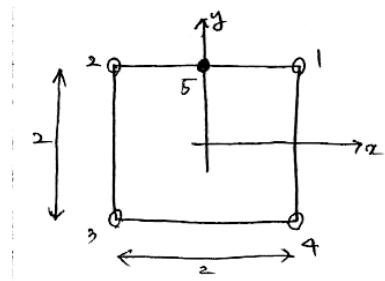
$$\mathbf{R}_S = \int_{-1}^{+1} \begin{bmatrix} \frac{1}{2}(1+x) & 0 \\ \frac{1}{2}(1-x) & 0 \\ 0 & \frac{1}{2}(1+x) \\ 0 & \frac{1}{2}(1-x) \end{bmatrix} \begin{bmatrix} 0 \\ -p(x) \end{bmatrix} \underbrace{(0.1)}_{\text{thickness}} dx \quad (6.12)$$

$$\mathbf{R}_S = \begin{bmatrix} 0 \\ 0 \\ -p_0(0.1) \\ -p_0(0.1) \end{bmatrix} \quad (6.13)$$



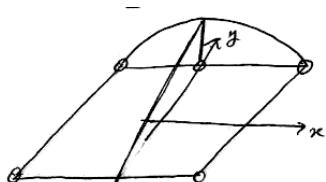
$$\text{total load} = P_0 \times 0.1 \times 2$$

6.1.2 Higher-order elements



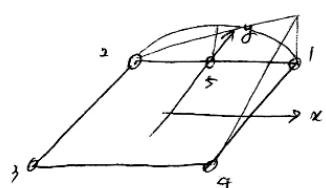
Want h_1, h_2, h_3, h_4, h_5

$$u(x, y) = \sum_{i=1}^5 h_i u_i.$$



$h_i = 1$ at node i and 0 at all other nodes.

$$h_5 = \frac{1}{2}(1-x^2)(1+y)$$



$$h_1 = \frac{1}{4}(1+x)(1+y) - \frac{1}{2}h_5 \quad (6.14)$$

$$h_2 = \frac{1}{4}(1-x)(1+y) - \frac{1}{2}h_5 \quad (6.15)$$

$$h_3 = \frac{1}{4}(1-x)(1-y) \quad (6.16)$$

$$h_4 = \frac{1}{4}(1+x)(1-y) \quad (6.17)$$

Note:

$$\boxed{\sum h_i = 1}$$

We must have $\sum_i h_i = 1$ to satisfy the rigid body mode condition.

$$u(x, y) = \sum_i h_i u_i \quad (6.18)$$

Assume all nodal point displacements = u^* . Then,

$$u(x, y) = \sum_i h_i u^* = u^* \sum_i h_i = u^* \quad (6.19)$$

From (6.1),

$$\left(\sum_m \mathbf{K}^{(m)} \right) \mathbf{U} = \mathbf{R} \quad (6.20)$$

$$\sum_m \left[\int_{V^{(m)}} \mathbf{B}^{(m)T} \mathbf{C}^{(m)} \mathbf{B}^{(m)} dV^{(m)} \right] \mathbf{U} = \mathbf{R} \quad (6.21)$$

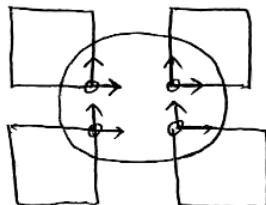
where $\mathbf{C}^{(m)} \mathbf{B}^{(m)} \mathbf{U} = \boldsymbol{\tau}^{(m)}$. (Assume we calculated \mathbf{U} .)

$$\sum_m \int_{V^{(m)}} \mathbf{B}^{(m)T} \boldsymbol{\tau}^{(m)} dV^{(m)} = \mathbf{R} \quad (6.22)$$

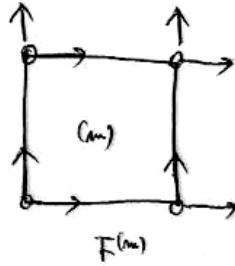
$$\sum_m \mathbf{F}^{(m)} = \mathbf{R}; \quad \mathbf{F}^{(m)} = \int_{V^{(m)}} \mathbf{B}^{(m)T} \boldsymbol{\tau}^{(m)} dV^{(m)} \quad (6.23)$$

Two properties

- I. The sum of the $\mathbf{F}^{(m)}$'s at any node is equal to the applied external forces.



II. Every element is in equilibrium under its $\mathbf{F}^{(m)}$



$$\hat{\mathbf{U}}^T \mathbf{F}^{(m)} = \underbrace{\hat{\mathbf{U}}^T \int_{V^{(m)}} \mathbf{B}^{(m)T} \boldsymbol{\tau}^{(m)} d V^{(m)}}_{= \bar{\epsilon}^{(m)T}} \quad (6.24)$$

$$= \int_{V^{(m)}} \bar{\epsilon}^{(m)T} \boldsymbol{\tau}^{(m)} d V^{(m)} \quad (6.25)$$

$$= 0 \quad (6.26)$$

where $\hat{\mathbf{U}}^T$ = virtual nodal point displacement.

Apply rigid body displacement.

If we move the element virtually in the rigid body modes, $\bar{\epsilon}^{(m)}$ is zero. Therefore the virtual work obtained due to virtual motion of the element is zero. Then the element is in equilibrium under its $\mathbf{F}^{(m)}$.