

MIT OpenCourseWare  
<http://ocw.mit.edu>

2.094 Finite Element Analysis of Solids and Fluids  
Spring 2008

For information about citing these materials or our Terms of Use, visit: <http://ocw.mit.edu/terms>.

Lecture 9 -  $u/p$  formulation

Prof. K.J. Bathe

MIT OpenCourseWare

We want to solve

Reading:  
Sec. 4.4.3

## I. Equilibrium

$$\begin{cases} \tau_{ij,j} + f_i^B = 0 & \text{in Volume} \\ \tau_{ij}n_j = f_i^{Sf} & \text{on } S_f \end{cases} \quad (9.1)$$

## II. Compatibility

## III. Stress-strain law

Use the principle of virtual displacements

$$\int_V \bar{\epsilon}^T C \epsilon \, dV = \mathcal{R} \quad (9.2)$$

We recognize that if  $\nu \rightarrow 0.5$ 

$$\epsilon_V \rightarrow 0 \quad (\epsilon_V = \epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz}) \quad (9.3)$$

$$\kappa = \frac{E}{3(1-2\nu)} \rightarrow \infty \quad (9.4)$$

$$p = -\kappa \epsilon_V \quad \text{must be accurately computed} \quad (9.5)$$

**Solution**

$$\tau_{ij} = \kappa \epsilon_V \delta_{ij} + 2G \epsilon'_{ij} \quad (9.6)$$

where

$$\delta_{ij} = \text{Kronecker delta} = \begin{cases} 1 & i = j \\ 0 & i \neq j \end{cases} \quad (9.7)$$

Deviatoric strains:

$$\epsilon'_{ij} = \epsilon_{ij} - \frac{\epsilon_V}{3} \delta_{ij} \quad (9.8)$$

$$\tau_{ij} = -p \delta_{ij} + 2G \epsilon'_{ij} \quad \left( p = -\frac{T_{kk}}{3} \right) \quad (9.9)$$

(9.2) becomes

$$\int_V \bar{\epsilon}^T C' \epsilon' \, dV + \int_V \bar{\epsilon}_V \kappa \epsilon_V \, dV = \mathcal{R} \quad (9.10)$$

$$\int_V \bar{\epsilon}^T C' \epsilon' \, dV - \int_V \bar{\epsilon}_V^T p \, dV = \mathcal{R} \quad (9.11)$$

We need another equation because we now have another unknown  $p$ .

$$p + \kappa \epsilon_V = 0 \quad (9.12)$$

$$\int_V \bar{p} (p + \kappa \epsilon_V) dV = 0 \quad (9.13)$$

$$-\int_V \bar{p} \left( \epsilon_V + \frac{p}{\kappa} \right) dV = 0 \quad (9.14)$$

For an element,

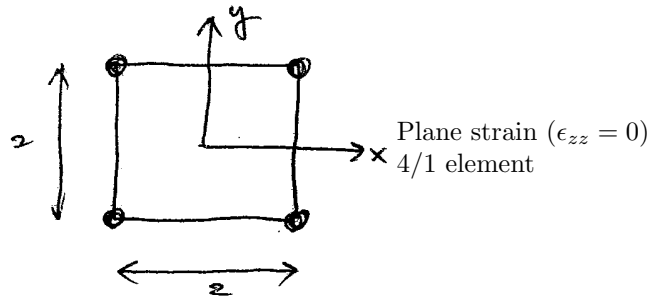
$$\mathbf{u} = \mathbf{H} \hat{\mathbf{u}} \quad (9.15)$$

$$\boldsymbol{\epsilon}' = \mathbf{B}_D \hat{\mathbf{u}} \quad (9.16)$$

$$\epsilon_V = \mathbf{B}_V \hat{\mathbf{u}} \quad (9.17)$$

$$p = \mathbf{H}_p \hat{p} \quad (9.18)$$

Example



Reading:  
Ex. 4.32 in  
the text

$$\epsilon_V = \epsilon_{xx} + \epsilon_{yy} \quad (9.19)$$

$$\boldsymbol{\epsilon}' = \begin{bmatrix} \epsilon_{xx} - \frac{1}{3}(\epsilon_{xx} + \epsilon_{yy}) \\ \epsilon_{yy} - \frac{1}{3}(\epsilon_{xx} + \epsilon_{yy}) \\ \gamma_{xy} \\ -\frac{1}{3}(\epsilon_{xx} + \epsilon_{yy}) \end{bmatrix} \quad (9.20)$$

**Note:**  $\epsilon_{zz} = 0$  but  $\epsilon'_{zz} \neq 0$ !

$$p = \mathbf{H}_p \hat{p} = [1] \{p_0\} \quad (9.21)$$

$$p(x, y) = p_0 \quad (9.22)$$

We obtain from (9.11) and (9.14)

$$\begin{bmatrix} \mathbf{K}_{uu} & \mathbf{K}_{up} \\ \mathbf{K}_{pu} & \mathbf{K}_{pp} \end{bmatrix} \begin{bmatrix} \hat{\mathbf{u}} \\ \hat{p} \end{bmatrix} = \begin{bmatrix} \mathbf{R} \\ \mathbf{0} \end{bmatrix} \quad (9.23)$$

$$\mathbf{K}_{uu} = \int_V \mathbf{B}_D^T \mathbf{C}' \mathbf{B}_D dV \quad (9.24a)$$

$$\mathbf{K}_{up} = - \int_V \mathbf{B}_V^T \mathbf{H}_p dV \quad (9.24b)$$

$$\mathbf{K}_{pu} = - \int_V \mathbf{H}_p^T \mathbf{B}_V dV \quad (9.24c)$$

$$\mathbf{K}_{pp} = - \int_V \mathbf{H}_p^T \frac{1}{\kappa} \mathbf{H}_p dV \quad (9.24d)$$

In practice, we use elements that use pressure interpolations per element, not continuous between elements. For example:



Then, unless  $\nu = 0.5$  (where  $\mathbf{K}_{pp} = \mathbf{0}$ ), we can use static condensation on the pressure dof's.

Use  $\hat{\mathbf{p}}$  equations to eliminate  $\hat{\mathbf{p}}$  from the  $\hat{\mathbf{u}}$  equations.

$$(\mathbf{K}_{uu} - \mathbf{K}_{up}\mathbf{K}_{pp}^{-1}\mathbf{K}_{pu})\hat{\mathbf{u}} = \mathbf{R} \quad (9.25)$$

(In practice,  $\nu$  can be 0.499999...)

The “best element” is the 9/3 element. (9 nodes for displacement and 3 pressure dof's).

$$p(x, y) = p_0 + p_1x + p_2y \quad (9.26)$$

### The inf-sup condition

Reading:  
Sec. 4.5

$$\underbrace{\inf}_{q_h \in Q_h} \underbrace{\sup}_{v_h \in V_h} \left[ \frac{\int_{\text{Vol}} q_h \overbrace{\nabla \cdot \mathbf{v}_h}^{=\epsilon_V} d\text{Vol}}{\underbrace{\|q_h\| \|v_h\|}_{\text{for normalization}}} \right] \geq \beta > 0 \quad (9.27)$$

$Q_h$ : pressure space.

If “this” holds, the element is optimal for the displacement assumption used (ellipticity must also be satisfied).

#### Note:

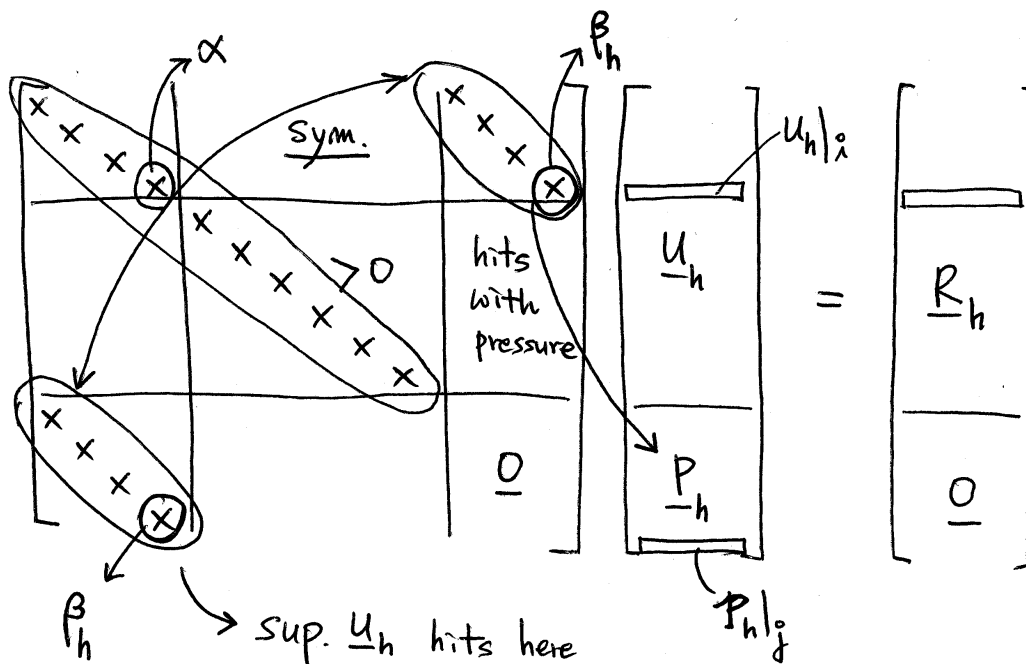
infimum = largest lower bound  
supremum = least upper bound

For example,

$$\begin{aligned} \inf \{1, 2, 4\} &= 1 \\ \sup \{1, 2, 4\} &= 4 \\ \inf \{x \in \mathbb{R}; 0 < x < 2\} &= 0 \\ \sup \{x \in \mathbb{R}; 0 < x < 2\} &= 2 \end{aligned}$$

(9.23) **rewritten** ( $\kappa = \infty$ , full incompressibility). Diagonalize using eigenvalues/eigenvectors.

For a mesh of element size  $h$  we want  $\boxed{\beta_h > 0}$  as we refine the mesh,  $h \rightarrow 0$



For  $\begin{pmatrix} x & x \\ x & x \end{pmatrix}$  (entry [3,1] in matrix) assume the circled entry is the minimum (inf) of  $\begin{pmatrix} x & x \\ x & x \end{pmatrix}$ .  
 Also, all entries in the matrix not shown are zero.

**Case 1**  $\beta_h = 0$

$$\Rightarrow \begin{cases} 0 \cdot u_h|_i = 0 & \text{(from the bottom equation)} \\ \underbrace{\alpha}_{\neq 0} \cdot u_h|_i + 0 \cdot p_h|_j = R_h|_i & \text{(from the top equation)} \end{cases}$$

$\Rightarrow$  no equation for  $p_h|_j$   
 $\Rightarrow$  spurious pressure! (any pressure satisfies equation)

**Case 2**  $\beta_h = \text{small} = \epsilon$

$$\epsilon \cdot u_h|_i = 0 \Rightarrow \boxed{u_h|_i = 0}$$

$$\therefore \epsilon \cdot p_h|_j + u_h|_i \cdot \alpha = R_h|_i$$

$$\Rightarrow p_h|_j = \frac{R_h|_i}{\epsilon} \Rightarrow \left( \begin{array}{l} \text{displ.} = 0 \\ \text{pressure} \rightarrow \text{large} \end{array} \right) \text{ as } \epsilon \text{ is small}$$

The behavior of given mesh when bulk modulus increases: locking, large pressures. See Example 4.39 textbook.